

## T Dependence of Equilibrium Constant

- **Basic issue**
  - we know  $K_p$  (or  $K_c$ ,  $K_f$ ), for a given reaction expression, is function of temperature only
  - can we describe/model the temperature dependence (for P.G.)
- **Approach**
  - employ previously derived expressions relating Gibbs Free Energy and Enthalpy (van't Hoff)

$$\left. \frac{\partial(G/T)}{\partial T} \right|_{p,n_i} = \frac{-H}{T^2} \quad \text{and for P.G.} \quad \frac{d(\mu_i^\circ/T)}{dT} = \frac{-\bar{h}_i}{T^2} = \frac{-\hat{h}_i}{T^2}$$

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## van't Hoff's Equation

- Start with  $K_p$  "definition" (P.G.)

$$\ln K_p = \frac{-1}{RT} \sum_i \nu_i \mu_i^\circ = \frac{-1}{R} \sum_i \nu_i \frac{\mu_i^\circ}{T}$$

- Get  $T$  dep. from derivative

$$\frac{d}{dT}(\ln K_p) = \frac{d}{dT} \left( \frac{-1}{R} \sum_i \nu_i \frac{\mu_i^\circ}{T} \right) = \frac{-1}{R} \sum_i \nu_i \frac{d}{dT} \left( \frac{\mu_i^\circ}{T} \right)$$

$$= \frac{1}{R} \sum_i \nu_i \frac{\hat{h}_i}{T^2}$$

$$\sum_i \nu_i \hat{h}_i = ?$$

$$= \hat{H}_{RHS}^\circ - \hat{H}_{LHS}^\circ \equiv \Delta \hat{H}_R^\circ \quad \text{Standard Heat of Reaction}$$

$$\Delta \hat{H}_R \begin{cases} < 0 & \text{Exothermic} \\ > 0 & \text{Endothermic} \end{cases}$$

drop  $^\circ$  since P.G.  $H \neq \hat{H}(p)$

$$\frac{d \ln K_p}{dT} = \frac{\Delta \hat{H}_R(T)}{RT^2} \quad \text{van't Hoff's Eq'n.}$$

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## $K_p$ Temperature Dependence

- So what is  $T$  dependence of  $K_p$

$$\frac{d}{dT}(\ln K_p) = \frac{\Delta \hat{H}_R(T)}{RT^2} \begin{cases} < 0 & \text{exothermic} \\ > 0 & \text{endothermic} \end{cases}$$

- From  $K_p$  def'n.  $K_p = \frac{\chi_{RHS}^{|\nu_i|}}{\chi_{LHS}^{|\nu_i|}} p^{\Delta n_R}$       $\Delta n_R \equiv \sum_i \nu_i$

– for exothermic reaction:  $T \uparrow \Rightarrow K_p \downarrow$  (more LHS)

– for endothermic reaction:  $T \uparrow \Rightarrow K_p \uparrow$  (more RHS)

- Recall  $p$  dependence

–  $p \uparrow \Rightarrow$  more lower moles side

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## $K_c$ Temperature Dependence

- For P.G., we had

$$\ln K_c = \ln K_p - \ln(\bar{R}T)^{\sum \nu_i} = \ln K_p - \Delta n_R \ln(\bar{R}T)$$

$$\begin{aligned} \frac{d}{dT} \ln K_c &= \frac{d}{dT} \ln K_p - \Delta n_R \frac{d}{dT} \ln(\bar{R}T) = \frac{\Delta \hat{H}_R}{RT^2} - \Delta n_R \frac{1}{T} \\ &= \frac{\Delta \hat{H}_R - \Delta n_R \bar{R}T}{RT^2} \end{aligned}$$

$$pV = n\bar{R}T \Rightarrow \Delta(pV) \underset{\text{fixed } T}{=} (\Delta n)\bar{R}T$$

$$= \frac{\Delta \hat{H}_R - \Delta pV}{RT^2}$$

$$\boxed{\frac{d}{dT} \ln K_c = \frac{\Delta \hat{U}_R}{RT^2}}$$

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## $K_f$ Temperature Dependence

- Using similar approach for imperfect gases

$$\frac{d}{dT} \ln K_f = \frac{\sum v_i \hat{h}_i^o}{RT^2}$$

- Pick  $f^o$  condition at low enough  $p$  such that  $f^o/p^o = 1$ 
  - we can choose any ref. ( $^o$ ) condition we want
  - by choosing low pressure  $\Rightarrow$  PG and  $\bar{h}_i = \hat{h}_i$   
and since  $K_f$  not function of  $f$

$$\frac{d}{dT} \ln K_f = \frac{\Delta H_R^o}{RT^2}$$

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## Interpolating for $K_p$

$$\frac{d \ln K_p}{dT} = \frac{\Delta \hat{H}_R(T)}{RT^2} \Rightarrow \int_{K_{p_1}}^{K_{p_2}} d \ln K_p = \ln(K_{p_2}/K_{p_1}) = \int_{T_1}^{T_2} \frac{\Delta \hat{H}_R(T)}{RT^2} dT$$

- Can usually find information on  $\Delta H_R$

$$\Delta \hat{H}_R = \sum v_i \hat{h}_i = \sum v_i \left( \hat{h}_{i,0} + \int_0^T \hat{c}_{p_i} dT \right) \cong \Delta \hat{H}_{R,0} + \sum v_i \int_0^T (a_i + b_i T + c_i T^2) dT$$

$$\cong \Delta \hat{H}_{R,0} + \sum v_i \left( a_i T + \frac{1}{2} b_i T^2 + \frac{1}{3} c_i T^3 \right)$$

- For “small”  $\Delta T$ , assume  $\Delta H_R \sim$  constant

$$\ln \frac{K_{p_2}}{K_{p_1}} \approx \frac{\Delta \hat{H}_R(T_1)}{R} \left[ \frac{1}{T_1} - \frac{1}{T_2} \right] \quad \text{i.e., 2<sup>nd</sup> term changes slowly works best for strongly exo. or endo. rxns.}$$

- Otherwise

$$\ln K_p \cong C + \frac{1}{R} \left\{ \frac{-\Delta \hat{H}_{R,T_0}}{T} + \sum v_i (a_i \ln T + b_i T/2 + c_i T^2/6) \right\}$$

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